

## Worksheet 8

### Problem 8.1: Separable-variables ODEs

- (a) Writing formally  $y'(x) = \frac{dy}{dx}$  and  $y(x) = y$ , we obtain

$$\frac{1}{y^2} dy = ax dx, \quad -\frac{1}{y} + C_1 = \frac{ax^2}{2} + C_2, \quad y(x) = \frac{2}{C - ax^2},$$

where  $C = 2(C_1 - C_2)$ . **Check:**

$$y'(x) = \frac{d}{dx} \frac{2}{C - ax^2} = -\frac{2 \cdot (-2ax)}{(C - ax^2)^2} = ax \frac{4}{(C - ax^2)^2} = ax y(x)^2.$$

- (b) Writing formally  $y'(x) = \frac{dy}{dx}$  and  $y(x) = y$ , we obtain

$$y dy = x^2 dx, \quad \frac{y^2}{2} + C_1 = \frac{x^3}{3} + C_2, \quad y(x) = \pm \sqrt{\frac{2}{3}x^3 + C},$$

where  $C = 2(C_2 - C_1)$ . **Check:**

$$y'(x) = \pm \frac{d}{dx} \left( \frac{2}{3}x^3 + C \right)^{1/2} = \pm \frac{1}{2} \left( \frac{2}{3}x^3 + C \right)^{-1/2} 2x^2 = \frac{x^2}{y(x)}.$$

- (c) This ODE seems to be **not a separable-variables** one. However, it is an (inhomogeneous) **linear ODE**,

$$y'(x) - y(x) = -x^2,$$

with inhomogeneity  $-x^2$ . One **particular** solution can be found by guessing,

$$y_{\text{part}}(x) = x^2 + 2x + 2.$$

Then, the **general** solution is obtained by adding to this **particular** solution the general solution of the corresponding **homogeneous** linear ODE  $y'(x) - y(x) = 0$ ,

$$y_{\text{gen}}(x) = (x^2 + 2x + 2) + Ae^x \quad (A \in \mathbb{R}).$$

- (d) Writing formally  $y'(x) = \frac{dy}{dx}$  and  $y(x) = y$ , we obtain

$$\frac{1}{1+y^2} dy = 4x^3 dx, \quad \arctan(y) + C_1 = x^4 + C_2, \quad y(x) = \tan(x^4 + C),$$

where  $C = C_2 - C_1$ , and we have used the first hint. **Check:**

$$y'(x) \equiv \frac{d}{dx} \tan(x^4 + C) = \frac{1}{\cos^2(x^4 + C)} 4x^3 = 4x^3 (1 + \tan^2(x^4 + C)).$$

- (e) Writing formally  $y'(x) = \frac{dy}{dx}$  and  $y(x) = y$ , we obtain

$$\cos y dy = 6x dx, \quad \sin y + C_1 = 3x^2 + C_2, \quad y(x) = \arcsin(3x^2 + C),$$

where  $C = C_2 - C_1$ . **Check:** Using the second hint, we get

$$y'(x) \equiv \frac{d}{dx} \arcsin(3x^2 + C) = \frac{1}{\sqrt{1 - (3x^2 + C)^2}} 6x.$$

Since  $y = \arcsin(3x^2 + C)$ , we have  $3x^2 + C = \sin y$ , and therefore

$$y'(x) = \frac{1}{\sqrt{1 - (\sin y)^2}} 6x = \frac{6x}{\cos y}.$$

(f) Writing formally  $y'(x) = \frac{dy}{dx}$  and  $y(x) = y$ , we obtain

$$\frac{dy}{\sin y} = dx.$$

We integrate the LHS, using the identity  $\sin y = 2 \sin \frac{y}{2} \cos \frac{y}{2}$ ,

$$\begin{aligned} \int dy \frac{1}{\sin y} &= \int dy \frac{1}{2 \sin \frac{y}{2} \cos \frac{y}{2}} = \int dy \frac{\frac{1}{2} \frac{1}{(\cos \frac{y}{2})^2}}{\tan \frac{y}{2}} \\ &= \int dy \frac{t'(y)}{t(y)} = \ln |t(y)| + C_1, \end{aligned}$$

where  $t(y) = \tan \frac{y}{2}$ , with the derivative  $t'(y) = \frac{1}{2} \frac{1}{(\cos \frac{y}{2})^2}$ .

Integrating the RHS,  $\int dx = x + C_2$  and setting  $C_2 - C_1 = C$ , we obtain

$$\ln \left| \tan \frac{y}{2} \right| = x + C.$$

Resolving for  $y$ , we eventually find<sup>1</sup>

$$y(x) = 2 \arctan (\pm e^{x+C}) = 2 \arctan (ae^x) \quad (a = \pm e^C \in \mathbb{R}).$$

**Check:** Since  $\frac{d}{du} \arctan u = \frac{1}{1+u^2}$ , the chain rule yields

$$\begin{aligned} y'(x) &= 2 \frac{ae^x}{1 + (ae^x)^2} = 2 \frac{\tan \frac{y(x)}{2}}{1 + \left(\tan \frac{y(x)}{2}\right)^2} \\ &= 2 \sin \frac{y(x)}{2} \cos \frac{y(x)}{2} \equiv \sin [y(x)]. \end{aligned}$$

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<sup>1</sup>Notice that  $\arctan(x)$  is the **inverse function** of  $\tan(y)$ :

$$\arctan(A) = B \Leftrightarrow \tan(B) = A.$$

### Problem 8.2

### Problem 8.3

- (a) This is a simple exercise in partial differentiation, see section 9.1.5.
- (b) This is a simple volume integral, see section 9.1.5.

### Problem 8.4

Using the ansatz  $T(x, y) = Ax^2 + Bxy + Cy^2 + Dx + Ey + F$  in the given PDE,

$$\nabla^2 T(x, y) \equiv \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) T = -\frac{s_0}{\lambda},$$

we find a condition on the values of  $A$  and  $C$ ,

$$2(A + C) = -\frac{s_0}{\lambda}. \quad (1)$$

Holding the rim of the disk at  $T = T_0$  corresponds to the (Dirichlet) boundary condition

$$T(x, y) = T_0 \quad (\text{for all } x, y \text{ with } x^2 + y^2 = R^2). \quad (2)$$

Setting  $B = D = E = 0$ , we see that conditions (1) and (2) can be satisfied by the choice

$$A = C = -\frac{s_0}{4\lambda}, \quad F = T_0 + \frac{s_0}{4\lambda} R^2 \quad (B = D = E = 0),$$

which turns the above ansatz into the (unique!) solution

$$T(x, y) = T_0 + \frac{s_0}{4\lambda} [R^2 - (x^2 + y^2)].$$

**Remark:** For completeness, we derive the given PDE (this is not part of Problem 8.4): In the medium, density  $\rho$  and current density  $\mathbf{J}$  of heat satisfy the **continuity equation**

$$\frac{\partial}{\partial t} \rho(\mathbf{r}, t) + \nabla \cdot \mathbf{J}(\mathbf{r}, t) = s(\mathbf{r}, t).$$

Using for  $\rho$  the **specific-heat relation**

$$\rho(\mathbf{r}, t) = c\mu [T(\mathbf{r}, t) - T_1]$$

(with some reference temperature  $T_1$ ), and for  $\mathbf{J}$  the **heat conduction law**,

$$\mathbf{J}(\mathbf{r}, t) = -\lambda \nabla T(\mathbf{r}, t),$$

we obtain a PDE for the temperature distribution  $T(\mathbf{r}, t)$  in the medium,

$$c\mu \frac{\partial}{\partial t} T(\mathbf{r}, t) - \lambda \nabla^2 T(\mathbf{r}, t) = s(\mathbf{r}, t).$$

In the case  $s(\mathbf{r}, t) = s(\mathbf{r})$ , steady-state equilibrium will be reached after some time, when  $\frac{\partial T}{\partial t} = 0$ . The resulting static distribution  $T(\mathbf{r}, t) = T(\mathbf{r})$  will satisfy the given PDE,

$$\nabla^2 T(\mathbf{r}) = -\frac{1}{\lambda} s(\mathbf{r}).$$

## Worksheet 9

### Problem 9.1: Finite difference method (FDM) for a PDE

- (a) In case of the **disk**  $\Sigma_1$ , the exact solution is obtained in Problem 4 of Worksheet 8, when we choose  $T_0 = 0$  and  $\frac{so}{\lambda} = S$  there,

$$T(x, y) = \frac{S}{4} [R^2 - (x^2 + y^2)].$$

(It is easy to check that this is really the solution !)

- (b) The five regions  $\Sigma_1, \dots, \Sigma_5$  in the  $xy$ -plane are shown in Fig. 1 ( $\Sigma_3$  is shown twice), along with a properly chosen mesh of discrete points in each case.

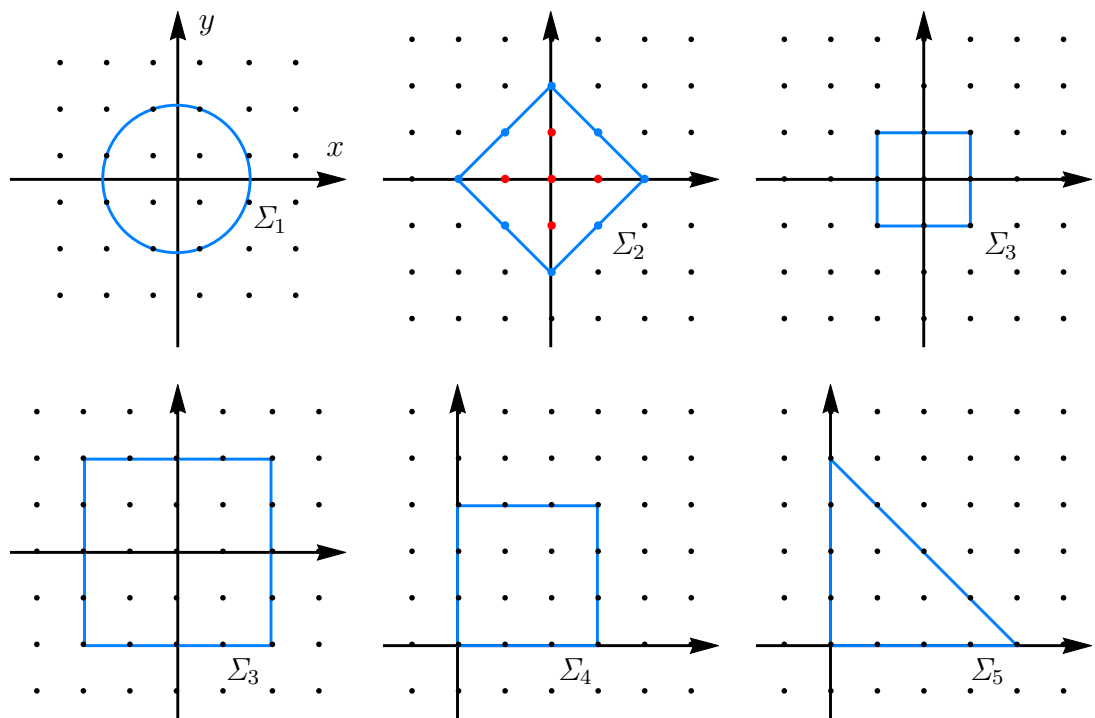


Figure 1: **Upper row:**  $\Sigma_1$  (with  $h = \frac{R}{5}\sqrt{10}$ ),  $\Sigma_2$  (with  $h = \frac{D}{4}$ ),  $\Sigma_3$  (with  $h = \frac{L}{2}$ ),  
**Lower row:**  $\Sigma_3$  (once more, but with  $h = \frac{L}{4}$ ),  $\Sigma_4$  (with  $h = \frac{L}{3}$ ), and  $\Sigma_5$ , (with  $h = \frac{L}{4}$ ).  
 The mesh points **inside**  $\Sigma_2$  are highlighted in red, the ones **on the rim**  $\partial\Sigma_2$  in blue.

**Remark 1:** Each panel in Fig. 1 has a certain number  $N$  of mesh points **inside**  $\Sigma_n$  (marked in red in the case of  $\Sigma_2$ ) plus a considerable number of extra mesh points **on the rim**  $\partial\Sigma_n$  (marked in blue in the case of  $\partial\Sigma_2$ ).

**Remark 2:** The four different panels with  $\Sigma_2, \Sigma_3, \Sigma_4$  all belong to the same physical situation (when we choose  $D = L\sqrt{2}$ ), but to different levels of approximation: With  $N = 5, 1, 9$ , or  $4$  mesh points inside  $\Sigma_n$  (the choice  $N = 1$  being a poor one).

- (c) In the lecture (section 12.2 "Finite difference methods"), we have obtained Eq. (448) ("a set of **iteration equations**") which in the present notation reads

$$T_{n,m} = \frac{1}{4} \left[ T_{n+1,m} + T_{n-1,m} + T_{n,m+1} + T_{n,m-1} \right] + \frac{h^2}{4} S_{n,m} .$$

In the present case  $S(x, y) = S$  is a constant,  $S_{n,m} = S$ , independently of  $n$  and  $m$ . In further simplified notation, we may write

$$T_{n,m} = \frac{T_{\rightarrow} + T_{\leftarrow} + T_{\uparrow} + T_{\downarrow}}{4} + \frac{h^2}{4} S , \quad (3)$$

where " $T_{\rightarrow}$ " means the temperature at the next-neighbor mesh point **to the right** (" $\rightarrow$ ") of  $\mathbf{r}_{n,m}$ , etc.

**Iteration method:** Using approximate values  $T_{\rightarrow}$ ,  $T_{\leftarrow}$ ,  $T_{\uparrow}$ ,  $T_{\downarrow}$  on the RHS of Eq. (3), the LHS should yield improved temperatures  $T_{n,m}$ .

**As an example,** we here pick the square  $\Sigma_2$  (second panel in upper row of Fig. 1): Given the point mesh of Fig. 1, we have 5 points inside  $\Sigma_2$  (marked in red) plus 8 ones on the rim  $\partial\Sigma_2$  (marked in blue). For short, we re-label the "red points"  $\mathbf{r}_{n,m}$  as follows

$$\begin{array}{cccc} & \mathbf{r}_{0,1} & & \mathbf{r}_2 \\ \mathbf{r}_{-1,0} & \mathbf{r}_{0,0} & \mathbf{r}_{1,0} & = & \mathbf{r}_3 & \mathbf{r}_5 & \mathbf{r}_1 \\ & \mathbf{r}_{0,-1} & & & & \mathbf{r}_4 & \end{array}$$

While the "blue points" are held at  $T = 0$ , we approximate the unknown temperatures  $T_n = T(\mathbf{r}_n)$  at the "red points" ( $n = 1, 2, 3, 4, 5$ ) by guessed values  $T_n^{(0)} \approx T_n$ . For example, choosing "heating strength"  $S$  and step size  $h$  as

$$S = 100, \quad h = 1,$$

a proper choice of these "initial guess" (or zeroth-order) temperatures  $T_n^{(0)}$  could be

$$\begin{array}{ccccccc} & & 0 & & & & 0 \\ & 0 & T_2^{(0)} & 0 & & 0 & 50 & 0 \\ 0 & T_3^{(0)} & T_5^{(0)} & T_1^{(0)} & 0 & = & 0 & 50 & 80 & 50 & 0 \\ & 0 & T_4^{(0)} & 0 & & & 0 & 50 & 0 \\ & & 0 & & & & & & 0 \end{array}$$

Using these trial values on the RHS of Eq. (3), we find the improved temperatures

$$\begin{aligned} T_1^{(1)} &= \frac{0 + 0 + T_5^{(0)} + 0}{4} + \frac{100}{4} = 45, \\ T_2^{(1)} &= \frac{0 + 0 + 0 + T_5^{(0)}}{4} + \frac{100}{4} = 45, \\ T_3^{(1)} &= \frac{T_5^{(0)} + 0 + 0 + 0}{4} + \frac{100}{4} = 45, \\ T_4^{(1)} &= \frac{0 + T_5^{(0)} + 0 + 0}{4} + \frac{100}{4} = 45, \\ T_5^{(1)} &= \frac{T_1^{(0)} + T_2^{(0)} + T_3^{(0)} + T_4^{(0)}}{4} + \frac{100}{4} = 75. \end{aligned}$$

These values, obtained upon the **first iteration**, comprise our **first-order result**,

$$\begin{array}{cccccc} & & 0 & & & 0 \\ & & 0 & T_2^{(1)} & 0 & & 0 & 45 & 0 \\ 0 & T_3^{(1)} & T_5^{(1)} & T_1^{(1)} & 0 & = & 0 & 45 & 75 & 45 & 0 \\ & & 0 & T_4^{(1)} & 0 & & & 0 & 45 & 0 \\ & & & 0 & & & & & 0 & & \end{array}$$

A **second iteration**, using these first-order temperatures  $T_n^{(1)}$  as new input on the RHS of Eq. (3), yields an even better set of **second-order temperatures**

$$T_1^{(2)} = T_2^{(2)} = T_3^{(2)} = T_4^{(2)} = 43.75, \quad T_5^{(2)} = 70.$$

One can (quite easily!) show that, after an infinite number of iterations, this procedure will converge towards the solution

$$T_1 = T_2 = T_3 = T_4 = \frac{5}{12} S \approx 41.67, \quad T_5 = \frac{2}{3} S \approx 66.67.$$

**Problem 9.2:**

- (a) This is an inhomogeneous linear ODE, with the general solution

$$\begin{aligned} f_{\text{gen}}^{\text{inh}}(x) &= f_{\text{gen}}^{\text{hom}}(x) + f_{\text{part}}^{\text{inh}}(x) \\ &= \left[ A \cos 7x + B \sin 7x \right] e^{-4x} - \frac{1}{5}, \end{aligned}$$

where  $f_{\text{gen}}^{\text{hom}}(x)$  can be found from an exponential ansatz  $f(x) = e^{\lambda x}$ ,

$$\lambda^2 + 8\lambda + 65 = 0, \quad \lambda_{1,2} = \frac{-8 \pm \sqrt{64 - 240}}{2} = \frac{-8 \pm 14i}{2} = -4 \pm 7i,$$

and  $f_{\text{part}}^{\text{inh}}(x) = -\frac{1}{5}$  (= const.) can be guessed.

- (b) This is another inhomogeneous linear ODE, with the general solution

$$\begin{aligned} f_{\text{gen}}^{\text{inh}}(x) &= f_{\text{gen}}^{\text{hom}}(x) + f_{\text{part}}^{\text{inh}}(x) \\ &= \left[ A + B e^{-8x} \right] - \frac{65}{8} x, \end{aligned}$$

where  $f_{\text{gen}}^{\text{hom}}(x)$  can be found from an exponential ansatz  $f(x) = e^{\lambda x}$ ,

$$\lambda^2 + 8\lambda = 0, \quad \lambda_{1,2} \in \{0, -8\},$$

and  $f_{\text{part}}^{\text{inh}}(x) = -\frac{65}{8} x$  can be guessed.

- (c)  $f''(x) + 8f'(x)^2 = 0$  is equivalent to a non-linear 1st-order ODE for  $g(x) = f'(x)$ ,

$$g'(x) = -8g(x)^2, \quad \frac{dy}{dx} = -8y^2, \quad -\frac{1}{y^2} dy = 8 dx, \quad \frac{1}{y} = 8x + C,$$

with the solution

$$\begin{aligned} g(x) &= \frac{1}{8x + C} \equiv f'(x) \\ \Rightarrow \quad f(x) &= \frac{1}{8} \ln(8x + C) + D \quad (C, D \in \mathbb{R}). \end{aligned}$$

### Problem 9.3: Average Temperature

(a) With  $V_\Omega = \frac{\pi}{2}R^2H$  and  $T(\mathbf{r}) \equiv T(x, y, z) = T_0 \frac{x}{R}$ , we have

$$\begin{aligned}
 \langle T(\mathbf{r}) \rangle_{\mathbf{r} \in \Omega} &= \frac{1}{V_\Omega} \int_\Omega d^3r T(\mathbf{r}) \\
 &= \frac{1}{\frac{\pi}{2}R^2H} \frac{T_0}{R} \int_0^H dz \int_{-R}^R dy \int_0^{\sqrt{R^2-y^2}} dx x \\
 &= \frac{2T_0}{\pi R^3 H} \underbrace{\int_0^H dz}_{=H} \int_{-R}^R dy \left[ \frac{x^2}{2} \right]_{x=0}^{x=\sqrt{R^2-y^2}} \\
 &= \frac{2T_0}{\pi R^3} \int_{-R}^R dy \left[ \frac{R^2 - y^2}{2} \right] \\
 &= \frac{2T_0}{\pi R^3} \cdot \frac{1}{2} \left[ R^2 y - \frac{y^3}{3} \right]_{y=-R}^{y=R} \\
 &= \frac{2T_0}{\pi R^3} \cdot \frac{1}{2} \left[ 2R^3 - \frac{2R^3}{3} \right] = \frac{4}{3\pi} T_0 \approx 0.424 T_0.
 \end{aligned}$$

(b) In cylindrical coordinates, with  $\tilde{T}(s, \phi, z) = T_0 \frac{s \cos \phi}{R}$  and the Jacobian  $J(s, \phi, z) = s$ ,

$$\begin{aligned}
 \langle T(\mathbf{r}) \rangle_{\mathbf{r} \in \Omega} &= \frac{1}{V_\Omega} \int_0^H dz \int_0^R ds \int_{-\pi/2}^{\pi/2} d\phi J(s, \phi, z) \tilde{T}(s, \phi, z) \\
 &= \frac{1}{\frac{\pi}{2}R^2H} \frac{T_0}{R} \underbrace{\int_0^H dz}_{=H} \int_0^R ds \int_{-\pi/2}^{\pi/2} d\phi s^2 \cos \phi \\
 &= \frac{2T_0}{\pi R^3} \left( \int_0^R ds s^2 \right) \left( \int_{-\pi/2}^{\pi/2} d\phi \cos \phi \right) = \frac{2T_0}{\pi R^3} \cdot \frac{R^3}{3} \cdot 2 = \frac{4}{3\pi} T_0.
 \end{aligned}$$

**Notice:** As the limits of each inner integral are fixed numbers, not depending on any of the respective outer integration variables, this triple integral **factorises**.