

Übungen zu Quantenfeldtheorie

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Blatt 2

Problem 1: Transition amplitude for a free particle

Show that the transition amplitude for a free particle with mass m moving in one dimension has the expression

$$\langle q', t' | q, t \rangle = \langle q' | e^{-i\hat{H}(t'-t)} | q \rangle = \left[\frac{m}{2\pi i(t' - t)} \right]^{1/2} \exp \left[\frac{im}{2} \frac{(q' - q)^2}{t' - t} \right] \quad (1)$$

To do the calculation, give t' a small imaginary part, $t' \rightarrow t' - i\epsilon$, and take $\epsilon \rightarrow 0^+$ at the end of the calculation. Show that this is justified, given that we interpret the position space propagator $\langle q', t' | q, t \rangle$ as a distribution, i.e. it is supposed to be integrated with an integrable wave function ψ .

$$\psi(q', t') = \int_{-\infty}^{\infty} dq \langle q', t' | q, t \rangle \psi(q, t). \quad (2)$$

Hint: Use the dominated convergence theorem.

Problem 2: Spreading of a wave packet

For the free particle, suppose $\psi(q, t = 0)$ is a Gaussian wave packet

$$\psi(q, t = 0) = \left(\frac{1}{2\pi\sigma_0^2} \right)^{1/4} \exp \left[-\frac{(q - a)^2}{4\sigma_0^2} \right].$$

Show that it will spread as time evolves:

$$|\psi(q, t)|^2 = \left(\frac{1}{2\pi\sigma(t)^2} \right)^{1/2} \exp \left[-\frac{(q - a)^2}{2\sigma(t)^2} \right],$$

where

$$\sigma(t) = \sqrt{\sigma_0^2 + \frac{t^2}{4m^2\sigma_0^2}}.$$

Hint: Use eqs. (1) and (2).

Problem 3: Non-standard path-integral representation

Consider the Lagrangian with a position-dependent “mass”

$$L = \frac{1}{2} m(q) \dot{q}^2.$$

Thus the canonical momentum is $p = m(q)\dot{q}$ and the Hamiltonian $H = \frac{1}{2m(q)}p^2$. Show that the path-integral representation of the transition amplitude has the form

$$\langle q' | e^{-i\hat{H}(t'-t)} | q \rangle = \mathcal{N} \int \mathcal{D}q \exp \left[i \int_t^{t'} ds \left(L(q(s), \dot{q}(s)) - \frac{i}{2} \tilde{\mathcal{N}} \log m(q(s)) \right) \right],$$

where $\tilde{\mathcal{N}} \equiv \lim_{\Delta t \rightarrow 0} \frac{1}{\Delta t}$ is an infinite constant, similar to \mathcal{N} .

Remark 1: Recall that the path integral is usually ill-defined in the continuum limit, where the discretization spacing Δt goes to zero. The continuum limit can only be taken for observable quantities.

Remark 2: This is a counterexample, given by Lee and Yang (1962), showing that the path-integral representation is not always of the “standard” form.

Problem 4: Poles in correlation functions

Consider a Fourier transform of the vacuum expectation value of a time-ordered product

$$T(q^2) = i \int d^4x e^{iqx} \langle \Omega | T \{ \hat{A}(x) \hat{A}^\dagger(0) \} | \Omega \rangle,$$

where the local operator $\hat{A}(x)$ can be either elementary scalar field operator like $\hat{\phi}(x)$ or composite operator like $\hat{\phi}^2(x)$ etc. The main restriction is that \hat{A} is composed of fields at the same space-time point, which makes it “local”.

Suppose \hat{A} has a non-zero matrix element between the vacuum and the one-particle state $|\vec{p}\rangle = \hat{a}_{\vec{p}}^\dagger |\Omega\rangle$. Here \vec{p} is the three-momentum and $p = (E(\vec{p}), \vec{p})$, where $E(\vec{p}) = \sqrt{m^2 + \vec{p}^2}$. Recall that in the relativistic normalization $\langle \vec{p} | \vec{q} \rangle = (2\pi)^3 2E(\vec{p}) \delta^{(3)}(\vec{p} - \vec{q})$.

Show that the function $T(q^2)$ considered as a function of the complex variable $q^2 = q_0^2 - \vec{q}^2$ has a pole in the vicinity of the point $q^2 = m^2$ and we can write

$$T(q^2) = \frac{|\langle \Omega | \hat{A}(0) | \vec{q} \rangle|^2}{m^2 - q^2 - i\epsilon}.$$

Check that this agrees with the expected result of the propagator in the non-interacting scalar field theory if $\hat{A} = \hat{\phi}$.

Hints:

- Write out the T product explicitly (in terms of Heaviside functions Θ) and insert

$$1_{\text{one-particle subspace}} = \int \frac{d^3p}{(2\pi)^3 2E(\vec{p})} |\vec{p}\rangle \langle \vec{p}|$$

in between the operators. Explain when inserting this “unity” is correct and when not. How does this completeness relation have to be modified to be valid more generally?

- Prove and use the relation $\langle \Omega | \hat{A}(x) | \vec{p} \rangle = \langle \Omega | \hat{A}(0) | \vec{p} \rangle e^{-ipx}$ (Hint: How does \hat{A} transform under translations?)
- Use the identity

$$\Theta(t) = \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} d\tau \frac{e^{it\tau}}{\tau - i\epsilon}.$$