

Quantum theory of condensed matter II

Mesoscopic physics

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Tue 8:00 - 10:00 9.2.01

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Fri 10:00 - 12:00 H33

Fri 12:00 - 14:00 5.0.20

Sheet 11

1. Markovian master equation for the Anderson impurity model

Let us consider an Anderson impurity coupled to an electronic lead as in Sheet 10

$$\hat{H} = \hat{H}_S + \hat{H}_B + \hat{H}_T$$

where

$$\hat{H}_S = \sum_{\sigma} \varepsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{\uparrow} \hat{n}_{\downarrow}, \quad (1a)$$

$$\hat{H}_B = \sum_{\mathbf{k}\sigma} \varepsilon_{\mathbf{k}} \hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma}, \quad (1b)$$

$$\hat{H}_T = \sum_{\mathbf{k}\sigma} \tau \left(\hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{d}_{\sigma} + \hat{d}_{\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma} \right). \quad (1c)$$

Assuming the high temperature limit ($k_B T \gg \hbar \gamma$ where $\gamma = \frac{2\pi\tau^2 D_0}{\hbar}$ and D_0 is the bath density of states at the Fermi level) we derived the following time local equation for the reduced density matrix, valid up to second order in the tunnelling Hamiltonian H_T ,

$$\begin{aligned} \dot{\hat{\rho}}_{\text{red}}(t) = & -\frac{\tau^2}{\hbar^2} \sum_{\sigma} \int_0^t dt' [F(t-t', +\mu) \hat{d}_{\sigma}(t) \hat{d}_{\sigma}^{\dagger}(t') \hat{\rho}_{\text{red}}(t') \\ & + F(t-t', -\mu) \hat{d}_{\sigma}^{\dagger}(t) \hat{d}_{\sigma}(t') \hat{\rho}_{\text{red}}(t') \\ & - F^*(t-t', -\mu) \hat{d}_{\sigma}(t) \hat{\rho}_{\text{red}}(t') \hat{d}_{\sigma}^{\dagger}(t') \\ & - F^*(t-t', +\mu) \hat{d}_{\sigma}^{\dagger}(t) \hat{\rho}_{\text{red}}(t') \hat{d}_{\sigma}(t') \\ & + \text{h.c.}], \end{aligned} \quad (2)$$

where all the operators are in the interaction picture and the bath correlation function is defined as

$$F(t-t', \mu) = \sum_{\mathbf{k}} \text{Tr}_B \left\{ \hat{c}_{\mathbf{k}\sigma}^{\dagger}(t) \hat{c}_{\mathbf{k}\sigma}(t') \hat{\rho}_B \right\}$$

1. Argue that, if we are interested into a time dynamics on time scales larger than the bath correlation time $\hbar\beta$, the time integration limit can be moved from the initial time $t_0 = 0$ to $t_0 = -\infty$ (Markov approximation).
2. Transform the equation from the interaction to the Schrödinger picture:

$$\begin{aligned} \dot{\hat{\rho}}_{\text{red}}(t) = & -\frac{i}{\hbar} \left[\hat{H}_S, \hat{\rho}_{\text{red}}(t) \right] - \frac{\tau^2}{\hbar^2} \sum_{\sigma} \int_0^{\infty} dt' [F(t', +\mu) \hat{d}_{\sigma} \hat{d}_{\sigma}^{\dagger}(-t') \hat{\rho}_{\text{red}}(t) \\ & + F(t', -\mu) \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma}(-t') \hat{\rho}_{\text{red}}(t) \\ & - F^*(t', -\mu) \hat{d}_{\sigma} \hat{\rho}_{\text{red}}(t) \hat{d}_{\sigma}^{\dagger}(-t') \\ & - F^*(t', +\mu) \hat{d}_{\sigma}^{\dagger} \hat{\rho}_{\text{red}}(t) \hat{d}_{\sigma}(-t') \\ & + \text{h.c.}]. \end{aligned} \quad (3)$$

where the density operators are in the Schrödinger picture, while the creation and annihilation operators of the impurity are still in the interaction picture.

3. Find the eigenenergies of the impurity system and write the equations for the populations in that basis using Eq.(3).
4. Considering the analytic expression of the correlator $F(t-t', \mu)$ we have found in Sheet 10, perform the time integral in Eq.(3) and obtain the master equation for the populations:

$$\dot{P}_0(t) = -2\gamma L(\varepsilon_d - \mu, W) f^+(\varepsilon_d) P_0(t) + \gamma L(\varepsilon_d - \mu, W) \sum_{\sigma} f^-(\varepsilon_d) P_{1\sigma}(t) \quad (4a)$$

$$\begin{aligned} \dot{P}_{1\sigma}(t) = & \gamma L(\varepsilon_d - \mu, W) f^+(\varepsilon_d) P_0(t) + \\ & -\gamma [L(\varepsilon_d + U - \mu, W) f^+(\varepsilon_d + U) + L(\varepsilon_d - \mu, W) f^-(\varepsilon_d)] P_{1\sigma}(t) + \\ & + \gamma L(\varepsilon_d + U - \mu, W) f^-(\varepsilon_d + U) P_2(t) \end{aligned} \quad (4b)$$

$$\dot{P}_2(t) = +\gamma \sum_{\sigma} L(\varepsilon_d + U - \mu, W) f^+(\varepsilon_d + U) P_{1\sigma}(t) - 2\gamma L(\varepsilon_d + U - \mu, W) f^-(\varepsilon_d + U) P_2(t) \quad (4c)$$

where $P_0(t) \equiv \langle 0 | \hat{\rho}_{\text{red}}(t) | 0 \rangle$, $P_{1\sigma} \equiv \langle 1\sigma | \hat{\rho}_{\text{red}}(t) | 1\sigma \rangle$ and $P_2(t) \equiv \langle 2 | \hat{\rho}_{\text{red}}(t) | 2 \rangle$ are the populations of the reduced density matrix with respect to the energy eigenbasis $|0\rangle, |1\uparrow\rangle, |1\downarrow\rangle, |2\rangle$ of the impurity. Moreover $f^+(\epsilon) \equiv [1 + \exp(\beta(\epsilon - \mu))]^{-1}$ and $f^-(\epsilon) \equiv f^+(-\epsilon)$.

In the stationary limit $\dot{P}_i = 0$ for $i \in \{ |0\rangle, |1\sigma\rangle, |2\rangle \}$. Is the linear system of equations well defined? What is the physical interpretation? How do we solve this issue?

Hint: Perform the integration with respect to the time difference $t-t'$ of the exponential dependence in $F(t-t', \mu)$ keeping into account that

$$\int_0^{+\infty} dx e^{-ax} e^{-ibx} = \frac{1}{a+ib} \quad \text{with } a > 0, b \in \mathbb{R}.$$

Moreover after the integration the following identity may be useful in order to sum the series

$$\sum_{k=0}^{\infty} \frac{x}{[(2k+1)\pi]^2 + x^2} \frac{y^2}{[(2k+1)\pi]^2} = \frac{1}{4} \left[\frac{y^2}{y^2 + x^2} \tanh\left(\frac{x}{2}\right) - \frac{xy}{y^2 + x^2} \tan\left(\frac{y}{2}\right) \right]$$

5. Prove that the stationary solution of the master equation is:

- i) $P_0 = 1, P_{1\sigma} = P_2 = 0$ for $\mu \ll \varepsilon_d$;
- ii) $P_2 = 1, P_{1\sigma} = P_0 = 0$ for $\mu \gg \varepsilon_d + U$;
- iii) $P_{1\sigma} = 1/2, P_2 = P_0 = 0$ for $\varepsilon_d \ll \mu \ll \varepsilon_d + U$;

where inequalities are taken with respect to the thermal energy $k_B T$ and the solution iii) is considered in the limit $U \gg k_B T$. Comment the result.

Frohes Schaffen!