

## Density Matrix Theory

Prof. Milena Grifoni  
Dr. Andrea Donarini

H33 Tuesdays, 10:15  
9.1.09 Tuesdays, 13:15

## Sheet 9

## 1. Real time dynamics of a q-bit coupled to an external bath

Consider the spin boson Hamiltonian:

$$H = H_S + H_{S-B} + H_B \quad (1)$$

where

$$\begin{aligned} H_S &= -\frac{\hbar}{2}(\epsilon\sigma_z + \Delta\sigma_x), \\ H_{S-B} &= \frac{1}{2}\sigma_z \sum_{\alpha} \hbar\nu_{\alpha}(b_{\alpha}^{\dagger} + b_{\alpha}) + \Delta V, \\ H_B &= \sum_{\alpha} \hbar\omega_{\alpha} b_{\alpha}^{\dagger} b_{\alpha}. \end{aligned}$$

The stochastic properties of  $H$  are characterized by the spectral density  $G(\omega) \equiv \sum_{\alpha} \nu_{\alpha}^2 \delta(\omega - \omega_{\alpha})$ .

1. Prove that the transformation to the energy eigenbasis of  $H_S$  can be written as a rotation

$$\begin{aligned} |1\rangle &= \cos \frac{\theta}{2} |\uparrow\rangle - \sin \frac{\theta}{2} |\downarrow\rangle \\ |2\rangle &= \sin \frac{\theta}{2} |\uparrow\rangle + \cos \frac{\theta}{2} |\downarrow\rangle \end{aligned}$$

where  $|1\rangle, |2\rangle$  are the energy eigenstates. Determine  $\theta$  as a function of  $\epsilon$  and  $\Delta$ .

2. Prove that the Hamiltonian (1), in the basis that diagonalizes  $H_S$  reads:

$$H = -\frac{\hbar}{2}\Delta_b\sigma'_z + \frac{1}{2}\left(\frac{\epsilon}{\Delta_b}\sigma'_z - \frac{\Delta}{\Delta_b}\sigma'_x\right) \sum_{\alpha} \hbar\nu_{\alpha}(b_{\alpha}^{\dagger} + d_{\alpha}) + \Delta V + H_B$$

where  $\Delta_b = \sqrt{\epsilon^2 + \Delta^2}$ .

3. Using the general expression for the time evolution of the two level system coupled to a bath derived in the lecture prove that (neglecting the Lamb shift):

$$\begin{aligned} \rho_{\uparrow\uparrow}(t) &= \cos \theta \left[ e^{-\Gamma_{rel}t} (\rho_{11}(0) - \rho_{11}(\infty)) + \rho_{11}(\infty) \right] + \sin^2 \frac{\theta}{2} \\ &\quad + \sin \theta e^{-\Gamma_{\phi}t} \left[ \text{Re}(\rho_{12}(0)) \cos(\Delta_b t) + \text{Im}(\rho_{12}(0)) \sin(\Delta_b t) \right] \end{aligned}$$

where  $\rho_{11}(\infty) = e^{-\beta E_1}/Z$  represent the thermodynamic limit and the relaxation and dephasing rates  $\Gamma_{rel}$  and  $\Gamma_{\phi}$  read respectively

$$\Gamma_{rel} = \frac{\pi}{2} \frac{\Delta^2}{\Delta_b^2} [G(-\Delta_b)N(-\Delta_b) + G(\Delta_b)N(\Delta_b)]$$

$$\Gamma_{\phi} = \frac{\Gamma_{rel}}{2} + \frac{\epsilon^2}{\Delta_b^2} \pi \lim_{\omega \rightarrow 0} G(\omega) \coth\left(\frac{\beta\hbar\omega}{2}\right)$$

where  $N(\omega) = \coth\left(\frac{\beta\hbar\omega}{2}\right) + 1$ .

4. The system-bath coupling is proportional to  $\sigma_z$ . The expectation value of the generalized dimensionless system coordinates reads correspondingly  $\langle \tilde{Q} \rangle = \rho_{\uparrow\uparrow} - \rho_{\downarrow\downarrow}$ . Calculate the time evolution of  $\langle \tilde{Q} \rangle$  for the initial condition  $\hat{\rho}(0) = |\uparrow\rangle\langle\uparrow|$ . Prove that the result is:

$$\langle \tilde{Q} \rangle(t) = \left[ \frac{\epsilon^2}{\Delta_b^2} - \langle \tilde{Q} \rangle(\infty) \right] e^{-\Gamma_{rel}t} + \langle \tilde{Q} \rangle(\infty) + \frac{\Delta^2}{\Delta_b^2} e^{-\Gamma_\phi t} \cos(\Delta_b t)$$

where  $\langle \tilde{Q} \rangle(\infty)$  is the expectation value of  $\tilde{Q}$  at equilibrium and can be obtained from the thermodynamic limit  $\rho_{ii} = \frac{e^{-\beta E_i}}{Z}$  where  $i = 1, 2$  label the energy eigenvalues of the system Hamiltonian.

## 2. Free propagator and classical action

Consider the propagator  $K(x_f t, x_i t_0)$  defined as by the relation:

$$K(x_f t, x_i t_0) \equiv \langle x_f | U(t, t_0) | x_i \rangle$$

as the matrix representation in space of the time evolver  $U(t, t_0)$  associated to the Schrödinger equation.

1. Prove that the propagator for the free particle can be written in the form:

$$K(x_f t, x_i t_0) = \sqrt{\frac{m}{2\pi\hbar i(t-t_0)}} \exp\left(\frac{i}{\hbar} \frac{m(x_f - x_i)^2}{2(t-t_0)}\right)$$

Hint: Use the fact that the Hamiltonian  $H$  for the free particle is time independent to get a simple expression of the propagator in terms of  $H$ . Simplify the expression of the propagator by making use of the completeness relation  $\mathbf{1} = \sum_n |\phi_n\rangle\langle\phi_n|$  where  $|\phi_n\rangle$  is a set of complete eigenstates of the Hamiltonian.

2. Prove that the quantum propagator introduced in the previous point can be written also in terms of the classical action:

$$K(x_f t, x_i t_0) = \sqrt{\frac{m}{2\pi\hbar i(t-t_0)}} \exp\left(\frac{i}{\hbar} S_{cl}\right)$$

where  $S_{cl}$  is the action functional of the classical path  $x_{cl}$ .

Hint: Start by proving that the classical path for a free particle can be parametrized with time as:

$$x_{cl}(s) = x_i + \frac{x_f - x_i}{t - t_0} (s - t_0)$$

and remember the definition of action  $S$  of the path  $x$

$$S[x] = \int_{t_0}^t ds \left( \frac{m\dot{x}^2}{2} - V(x) \right)$$

**Frohes Schaffen!**