

## Density Matrix Theory

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H33 Tuesdays, 10:15  
9.1.09 Tuesdays, 13:15

## Sheet 6

## 1. Equilibrium: the free energy formulation

Consider the master equation for the Anderson impurity model introduced in the Sheet 4:

$$\begin{aligned}\dot{P}_0 &= -2\gamma f^+(\varepsilon_d)P_0 + \gamma \sum_{\sigma} f^-(\varepsilon_d)P_{1\sigma} \\ \dot{P}_{1\sigma} &= -\gamma[f^+(\varepsilon_d + U) + f^-(\varepsilon_d)]P_{1\sigma} \\ &\quad + \gamma f^+(\varepsilon_d)P_0 + \gamma f^-(\varepsilon_d + U)P_2 \\ \dot{P}_2 &= -2\gamma f^-(\varepsilon_d + U)P_2 + \gamma \sum_{\sigma} f^+(\varepsilon_d + U)P_{1\sigma}\end{aligned}$$

where

$$P_0(t) \equiv \langle 0 | \rho_{red}(t) | 0 \rangle P_{1\sigma} \equiv \langle 1\sigma | \rho_{red}(t) | 1\sigma \rangle P_2(t) \equiv \langle 2 | \rho_{red}(t) | 2 \rangle$$

are the populations of the reduced density matrix with respect to the manybody energy eigenbasis  $|0\rangle$ ,  $|1\uparrow\rangle$ ,  $|1\downarrow\rangle$ ,  $|2\rangle$  of the impurity.

1. Prove that the stationary solution of this master equation is independent of the magnitude of the bare tunnelling rate  $\gamma$  and, for every value of the parameters  $(\varepsilon_d, U, \mu, T)$  defining the model, can be written in the form:

$$\begin{aligned}P_0^{stat} &= \frac{1}{N} f^-(\varepsilon_d) f^-(\varepsilon_d + U) \\ P_{1\sigma}^{stat} &= \frac{1}{N} f^+(\varepsilon_d) f^-(\varepsilon_d + U) \\ P_2^{stat} &= \frac{1}{N} f^+(\varepsilon_d) f^+(\varepsilon_d + U)\end{aligned}\tag{1}$$

where  $N$  is the normalization factor that ensures the sum of the probability to be 1. Moreover  $f^+(\varepsilon) \equiv [1 + e^{\beta(\varepsilon - \mu)}]^{-1}$  and  $f^-(\varepsilon) \equiv 1 - f^+(\varepsilon)$ .

2. Prove that the equilibrium probabilities derived at the previous point can be obtained from a thermodynamical formulation of the problem where the impurity, defined by the Hamiltonian  $H_S$  (see Sheet 4), can exchange energy and particles with a bath with temperature  $T$  and chemical potential  $\mu$ . In particular calculate the grandcanonical partition function  $\mathcal{Z} = \text{Tr}_S \{ e^{-\beta(H_S - \mu N_S)} \}$  for the impurity and prove that:

$$P_{\alpha}^{stat} = \frac{1}{\mathcal{Z}} \text{Tr}_S \{ |\alpha\rangle \langle \alpha| e^{-\beta(H_S - \mu N_S)} \}$$

where  $|\alpha\rangle$  is a manybody energy eigenstate of the impurity and  $N_S$  the particle number.

## 2. Time evolution for a Markovian master equation

In this exercise we consider the Markoff master equation (1) and calculate numerically the time evolution for the populations of the many-body states of the impurity.

1. Show that the equations (1) can be cast into a matrix form  $\dot{P}(t) = LP(t)$  where  $P \equiv (P_0, P_{1\uparrow}, P_{1\downarrow}, P_2)^T$  and

$$L = \gamma \begin{pmatrix} -2f^+(\varepsilon_d) & f^-(\varepsilon_d) & f^-(\varepsilon_d) & 0 \\ f^+(\varepsilon_d) & -f^-(\varepsilon_d) - f^+(\varepsilon_d + U) & 0 & f^-(\varepsilon_d + U) \\ f^+(\varepsilon_d) & 0 & -f^-(\varepsilon_d) - f^+(\varepsilon_d + U) & f^-(\varepsilon_d + U) \\ 0 & f^+(\varepsilon_d + U) & f^+(\varepsilon_d + U) & -2f^-(\varepsilon_d + U) \end{pmatrix}.$$

Prove that the solution of the equation can be written in the form  $P(t) = e^{Lt}P(t=0)$ . Taking advantage of this algebraic formulation, calculate the numerical solution of (1). Hint: the function “expm” calculates the exponential of a matrix in Matlab.

2. Prove that, if the time is measured in units of  $1/\gamma$  solutions with different tunneling rates coincide.
3. Check that the stationary solution is reached by the system after a time corresponding to a few  $1/\gamma$  and that it is independent of the initial condition.
4. Calculate the time evolution for the population vector  $P$  also with the help of one of the packages for ordinary differential equations available in Matlab. Compare the results with the previous method. Hint: There are different types of solvers. You can start by typing “help ode23” in the command line and read the documentation.

**Frohes Schaffen!**