

## Mesoscopic Physics

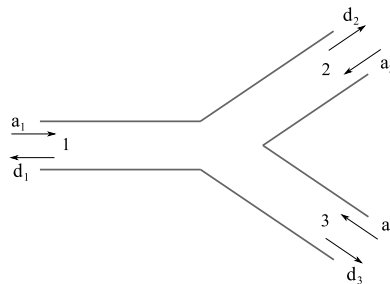
Dr. Andrea Donarini  
 Jürgen Wurm  
 Matthias Scheid

Room 3.1.26  
 Fridays at 10:15

### Sheet 10

#### 1. Mesoscopic beam splitter

Consider a device with three terminals, up/down symmetry and time reversal symmetry.



- (a) • Show that the scattering matrix can be parametrized as

$$S = \begin{pmatrix} r_0 & t & t \\ t & r & r' \\ t & r' & r \end{pmatrix} \quad (1)$$

- (b) • Assume real parameters and show that for nonzero  $t$  either

$$t^2 = \frac{1 - r_0^2}{2}, \quad r = -\frac{1 + r_0}{2}, \quad r' = \frac{1 - r_0}{2} \quad (2)$$

or

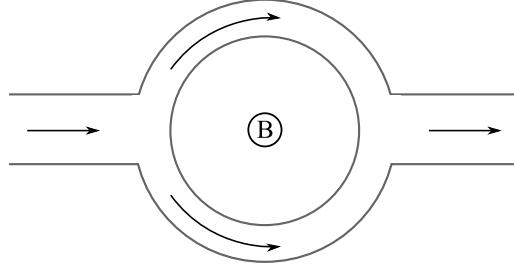
$$t^2 = \frac{1 - r_0^2}{2}, \quad r = \frac{1 - r_0}{2}, \quad r' = -\frac{1 + r_0}{2} \quad (3)$$

has to hold. What is the maximum value for  $t^2$ ?

- (c) • Consider a fully symmetric system. What changes? Can  $r$  become zero?

## 2. Mesoscopic Aharonov-Bohm effect

The conductance through a loop which is pierced by a magnetic field  $B$  oscillates as a function of the field. Consider two identical beam splitters with scattering matrices as in equation (1) connected in series through a ring threaded by a magnetic field. Assume that the magnetic field is nonzero only in the middle region of the ring and the electrons do not feel a Lorentz force.



- (a) Start with  $B = 0$ . Suppose that electrons acquire a phase  $\varphi$  when traversing either the upper or the lower branch of the ring (that is to say the “transmission” amplitude of one branch is  $e^{i\varphi}$  just as it is  $e^{ikL}$  for propagation through a piece of free space with length  $L$ ). Show that the total transmission amplitude  $\tilde{t}$  is given by

$$\tilde{t}(\varphi) = 2t^2 e^{i\varphi} \frac{1 - (r - r')^2 e^{2i\varphi}}{1 - 2(r^2 + r'^2) e^{2i\varphi} + (r^2 - r'^2)^2 e^{4i\varphi}}. \quad (4)$$

You can use Maple or Mathematica for the algebra. Show that for real parameters, using the results of problem 1, that

$$T \equiv |\tilde{t}|^2 = \frac{(1 - r_0^2)^2}{1 - 2r_0^2 \cos(2\varphi) + r_0^4}. \quad (5)$$

Where are the conductance resonances as a function of  $\varphi$ ? What does the resonance condition mean?

- (b) If the magnetic field is finite, an electron moving clockwise through one of the arms acquires an additional phase  $\phi$ , while an electron moving counterclockwise acquires an additional phase  $-\phi$  with  $2\phi = \oint \mathbf{A} \cdot d\mathbf{l} = 2\pi\Phi/\Phi_0$ . Here  $\Phi$  is the magnetic flux through the ring and  $\Phi_0 = h/e$  is the magnetic flux quantum. Show that in this case one gets

$$\tilde{t}(\varphi, \phi) = 2t^2 \cos(\phi) e^{i\varphi} \frac{1 - (r - r')^2 e^{2i\varphi}}{1 - 2(r^2 + r'^2 \cos[2\phi]) e^{2i\varphi} + (r^2 - r'^2)^2 e^{4i\varphi}}. \quad (6)$$

The transmission probability  $|\tilde{t}(\varphi, \phi)|^2$  is an oscillating function of  $\Phi$ . The fundamental frequency is given by  $\Phi_0$ . These oscillations are called Aharonov-Bohm (AB) oscillations. However also higher harmonics with periods that are integer fractions of  $\Phi_0$  are present, for example the Altshuler-Aronov-Spivak (AAS) oscillations with period  $\Phi_0/2$ . What is the physical origin of the AB and the AAS oscillations? Plot  $T(\varphi, \phi)$  for different parameters.

- (c) • Consider the limit of a nearly closed ring  $r_0 = 1 - \Delta$  and  $r' = \Delta/2$ ,  $\Delta \ll 1$ . Show that only the fundamental oscillation survives in leading order in  $\Delta$ . Explain this observation.

## 3. Remember the numerical problem of last week!

**Frohes Schaffen!**