Quantum Theory of Condensed Matter

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Sheet 2

1. Bosonic commutation relations

Refresh the physics of the simple harmonic oscillator

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{m\omega^2 \hat{x}^2}{2},$$

which can be written in "second quantized" form, by expressing \hat{x} and \hat{p} in terms of *boson* creation and annihilation operators:

$$\hat{H} = \hbar\omega \left(b^{\dagger}b + \frac{1}{2} \right), \quad b^{\dagger} = \frac{1}{\sqrt{2\hbar}} \left(\sqrt{m\omega} \, \hat{x} - \mathrm{i} \frac{\hat{p}}{\sqrt{m\omega}} \right).$$

From the canonical commutation relations between position and momentum operators, it follows immediately (do you remember it?) that the basis commutation relations hold:

$$[b, b^{\dagger}] = 1, \quad [b, b] = 0, \quad b |0\rangle = 0$$

where [A, B] = AB - BA, $|0\rangle$ is the vacuum, and \dagger indicates the Hilbert space adjoint.

a) Show that for two non commuting bosonic operators A, and B it holds

$$[A, B^{n}] = \sum_{k=0}^{n-1} B^{k} [A, B] B^{n-1-k}.$$

b) Prove, using the result of point a), the following relations valid for bosonic operators b, b^{\dagger}

$$\begin{bmatrix} b, (b^{\dagger})^n \end{bmatrix} = n(b^{\dagger})^{n-1} = \frac{\partial (b^{\dagger})^n}{\partial b^{\dagger}}, \\ \begin{bmatrix} b^{\dagger}, b^n \end{bmatrix} = -nb^{n-1} = -\frac{\partial b^n}{\partial b}.$$

2. Classical vibration in the diatomic linear crystal

Consider a periodic chain of atoms with alternating masses m, and M connected by springs with spring constant K. The equilibrium position of the atoms is na with $n \in \mathbb{Z}$.

a) Prove that the dispersion relation of the eigenmodes reads:

$$\omega(q)^{2} = \frac{K}{Mm} \left[M + m \pm \sqrt{M^{2} + m^{2} + 2mM\cos(2qa)} \right]$$

- b) Make a sketch of the dispersion relation and discuss the dynamics of the two modes in the limiting cases $qa \rightarrow 0$ and $2qa \rightarrow \pi$.
- c) Show that, in the limit m = M the dispersion reduces to the one of monoatomic linear chain.