

Assignments to Condensed Matter Theory I

Sheet 10

G. Cuniberti (Phy 4.1.29)
A. Donarini (Phy 3.1.24)

Room H35
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sheet online: <http://www-MCG.uni-R.de/teaching/>

Problem set: Electron-Electron Interaction

10.1. Occupation number representation

Let us consider a fermionic system with two single particle states $|\phi_1\rangle$ and $|\phi_2\rangle$ that span the (two-dimensional) single particle Hilbert space.

- (a) Which dimension has the two-particle Hilbert space? Which dimension has the Fock space? Write down the form of the basis of the Fock space explicitly as Slater determinants of the states ϕ_1 , ϕ_2 and in the occupation number representation.
- (b) Calculate in this basis the matrix representation of the creation and annihilation operators c_i, c_i^\dagger ($i = 1, 2$) and also of the occupation operators $n_i = c_i^\dagger c_i$
- (c) Calculate the anticommutator relations

$$[c_i, c_j]_+ = [c_i^\dagger, c_j^\dagger]_+ = 0; \quad [c_i, c_j^\dagger]_+ = \delta_{ij}$$

explicitly using matrix multiplication of the matrices calculated at point (b).

- (d) Consider an Hamilton operator

$$H = T + V$$

where T is a single particle operator and V a two particle one. With respect to the single particle basis $|\phi_i\rangle$ the matrix elements are:

$$\langle \phi_i | T | \phi_i \rangle = \epsilon; \quad \langle \phi_i | T | \phi_j \rangle = t \text{ for } i \neq j$$

$${}^{(1)}\langle \phi_1 | {}^{(2)}\langle \phi_2 | V | \phi_2 \rangle {}^{(2)} | \phi_1 \rangle {}^{(1)} = U; \quad {}^{(1)}\langle \phi_1 | {}^{(2)}\langle \phi_2 | V | \phi_1 \rangle {}^{(2)} | \phi_2 \rangle {}^{(1)} = J$$

where the notation is such that, *e.g.*:

$${}^{(1)}\langle \phi_1 | {}^{(2)}\langle \phi_2 | V | \phi_1 \rangle {}^{(2)} | \phi_2 \rangle {}^{(1)} \equiv \frac{1}{V^2} \int d\mathbf{r}_1 d\mathbf{r}_2 \phi_1^*(\mathbf{r}_1) \phi_2^*(\mathbf{r}_2) V(\mathbf{r}_1, \mathbf{r}_2) \phi_1(\mathbf{r}_2) \phi_2(\mathbf{r}_1)$$

Write the operator H in second quantization and in the matrix representation (starting from the single particle basis introduced). Calculate the eigenvalues and eigenvectors for H .

- (e) Again, write H in second quantization, but this time as a single particle basis use the eigenvectors of T . Which is the connection between this creation and annihilation operators and the ones considered in the points (a)-(d)? Is this a unitary transformation?

10.2. Wick's theorem

- (a) Show that, for a system of non-interacting fermions described by the Hamiltonian in the eigenvalue basis

$$H = \sum_{\alpha} \epsilon_{\alpha} c_{\alpha}^{\dagger} c_{\alpha},$$

the following relation for the many-body grandcanonical expectation values holds:

$$\langle c_{\alpha_1}^{\dagger} c_{\alpha_2}^{\dagger} c_{\alpha_3} c_{\alpha_4} \rangle = \langle c_{\alpha_1}^{\dagger} c_{\alpha_4} \rangle \langle c_{\alpha_2}^{\dagger} c_{\alpha_3} \rangle \delta_{\alpha_1 \alpha_4} \delta_{\alpha_2 \alpha_3} - \langle c_{\alpha_1}^{\dagger} c_{\alpha_3} \rangle \langle c_{\alpha_2}^{\dagger} c_{\alpha_4} \rangle \delta_{\alpha_1 \alpha_3} \delta_{\alpha_2 \alpha_4},$$

where

$$\langle c_{\alpha_1}^{\dagger} c_{\alpha_2}^{\dagger} c_{\alpha_3} c_{\alpha_4} \rangle \equiv \frac{1}{Z} \text{Tr} \{ c_{\alpha_1}^{\dagger} c_{\alpha_2}^{\dagger} c_{\alpha_3} c_{\alpha_4} \exp[-\beta(H - \mu N)] \}$$

and Z is the grandcanonical partition function. The trace is taken over the full Fock space.

- (b) Derive from point (a) that, for non-interacting fermions, in every other given single particle basis $\{|n\rangle\}$ the following relation holds:

$$\langle c_{n_1}^{\dagger} c_{n_2}^{\dagger} c_{n_3} c_{n_4} \rangle = \langle c_{n_1}^{\dagger} c_{n_4} \rangle \langle c_{n_2}^{\dagger} c_{n_3} \rangle - \langle c_{n_1}^{\dagger} c_{n_3} \rangle \langle c_{n_2}^{\dagger} c_{n_4} \rangle.$$

Note that this is valid even if in this basis the Hamiltonian

$$H = \sum_{n,m} t_{nm} c_n^{\dagger} c_m$$

would contain non-diagonal terms t_{nm} for $n \neq m$.

Hint: Diagonalize first H using a unitary transformation $c_n = \sum_{\alpha} u_{n\alpha} c_{\alpha}$. Apply then the equation proved in point (a). Finally perform the canonical transformation in the opposite direction.

10.3. [Kür] Model of two interacting particles in 1D

Let us consider two interacting particles to model a helium atom in 1D. In properly chosen dimensionless coordinates, write the first quantization Hamiltonian as:

$$H = -\frac{d^2}{dx_1^2} - 2V(|x_1|) - \frac{d^2}{dx_2^2} - 2V(|x_2|) + V(|x_1 - x_2|)$$

where

$$V(x) = \frac{2}{x + \delta}$$

is the "truncated" one dimensional Coulomb potential. The factor of two in the single particle potential is due to the double positive charge of the "Helium" nucleus.

- (a) Calculate the Hamiltonian in second quantization in the form:

$$H = \sum_{i,\sigma} \epsilon_i c_{i\sigma}^\dagger c_{i\sigma} + \frac{1}{2} \sum_{i_1 i_2 i_3 i_4, \sigma \sigma'} u_{i_1 i_2 i_3 i_4} c_{i_1 \sigma}^\dagger c_{i_2 \sigma'}^\dagger c_{i_3 \sigma'} c_{i_4 \sigma}$$

relative to the basis $|i\rangle$ of the eigenvectors of the single particle Hamiltonian. The single particle eigenfunctions $\phi_i(x) = \langle x|i\rangle$ fulfill the Schrödinger equation:

$$\left(-\frac{d^2}{dx^2} - 2V(|x|) \right) \phi_i(x) = \epsilon_i \phi_i(x).$$

In other words you have to calculate the eigenvalues ϵ_i and the Coulomb matrix elements u for $i, i_l \in \{1, 2\}$ numerically.

Hint: For the numerical calculation discretize the space $x \rightarrow x_n$ with $n = 1, \dots, N$. Now the wave-function is a vector since $\phi(x) \rightarrow \phi(x_n)$. The discrete version of the derivative now reads $\phi'(x_n) = (\phi(x_{n+1}) - \phi(x_{n-1})) / (x_{n+1} - x_{n-1})$. Work out analogously the second derivative and, finally, remember that the potential operator acts locally $(V\phi)(x) = V(x)\phi(x)$. Put all together and you have transformed the single particle Schrödinger equation into an algebraic equation that can be solved numerically.

- (b) Calculate the lowest energy two-particle eigenstates exactly and in Hartree-Fock approximation under the assumption that you can consider only the lowest single particle quantum number (that is $i, i_l \in \{1, 2\}$). Treat separately the singlet- (antiparallel spins, “ortho-helium”) and the triplet-case (total spin 1, “para-helium”).